Lecture 14

4. Hamiltonian Dynamics

Idea: Reformulate a dynamical system described by a Lagrangian I in terms of $q_i, P_i = \frac{\partial \mathcal{L}}{\partial \dot{q}_i}, i \in \{1, ..., N\}$

solve this for qi in terms of qi, pi

Start with the Lagrangian descriptions - start with co-ordinates $q_1, \dots, q_N \equiv q_N$ - set of conjugate momenta $p_1, \ldots, p_n \equiv p$ where $P_{i} = \frac{\partial \mathcal{L}}{\partial \dot{q}_{i}}$

Step 1: Define q_i in terms of p, q by solving relation for q_i

$$P_i = \frac{\partial L}{\partial \dot{a}}$$
i $\in \{1, ..., N\}$

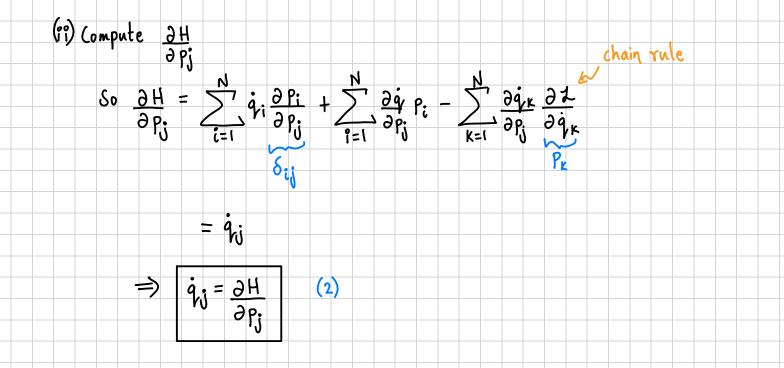
<u>Step 2</u>: Define Hamiltonian $H(q, p) = \sum_{i=1}^{N} \dot{q_i}(q, p)$ $H(q, p) = \sum_{i=1}^{N} \dot{q_i}(q, p)$

(i) Compute $\frac{\partial H}{\partial q_{j}}$ $i \in \{1, ..., N\}$ $\frac{\partial H}{\partial q_{j}} = \sum_{i=1}^{N} \frac{\partial q_{i}}{\partial q_{i}} p_{i} - \frac{\partial z}{\partial q_{i}} - \sum_{k=1}^{N} \frac{\partial q_{k}}{\partial q_{i}} \frac{\partial z}{\partial q_{k}}$

$$= -\frac{\partial d}{\partial q_{i}} = \frac{d}{dt} \frac{\partial d}{\partial q_{i}} = -\dot{p}_{i}$$

1

and Hherefore $\dot{p}_{j} = -\frac{\partial H}{\partial q_{j}}$ (1)



(1) and (2) together are Hamilton's eqn for the system

- set of 2N first order equations for the time evolution of q, p a set of 2N variables

- set of p and q form the phase space for the system

4.2 Examples

(1) Particle of mass m on the x axis subjected $\xrightarrow[]{0, x} \longrightarrow x$ to a force $\xrightarrow[]{0, x} \longrightarrow x$

$$F = -dV$$

$$\int_{0} f(x, \dot{x}) = \lim_{n \to \infty} \dot{x}^{2} - V(x)$$

The momentum conjugate to x is

$$p = \frac{\partial J}{\partial \dot{x}} = m\dot{x} \implies \dot{x} = \frac{P_x}{m}$$

Then the Hamiltonian is

$$\dot{x}_{p} - \lambda = \frac{p^{2}}{m} - \left(\frac{1}{2}mp^{2} - V\right)$$
$$\implies H = \frac{p^{2}}{m} + V$$
$$= \frac{p^{2}}{2m}$$

Then Hamilton's equations are

$$\dot{X} = \frac{\partial H}{\partial P_{x}} = \frac{P}{m}$$
Nautons equation

$$dP = -dV = mi = -dV$$

$$\dot{P} = -\frac{\partial H}{\partial x} = -\frac{\partial V}{\partial x}$$
(2) Particle on a plane
Use polar co-ordinates subject to a force

$$E = -\overline{Q}(V(\tau)) = -\frac{\gamma}{T}V'(\tau)$$

$$K.E = \frac{1}{2}m(\dot{\tau}^{2} + \tau^{2}\dot{\theta}^{2}), PE = V(\tau), \Delta = KE - PE$$

$$\Rightarrow \Delta(Y, \theta, \dot{\tau}, \dot{\theta}) = \frac{m}{2}(\dot{\tau}^{2} + \tau^{2}\dot{\theta}^{2}) - V(\tau)$$
and $P_{T} = \frac{\partial A}{\partial \dot{\tau}} = m\dot{\tau} \Rightarrow \dot{\tau} = \frac{P_{v}}{m}$
Further,
 $P_{\theta} = \frac{\partial A}{\partial \dot{\theta}} = m\tau^{2}\dot{\theta} = \dot{\theta} = \frac{P_{\theta}}{m\tau^{1}}$
The Hamiltonian is

$$H(\tau_{1}\theta, \dot{\tau}, \dot{\theta}) = \dot{\tau}P_{r} + \frac{P_{\theta}}{2m} - (\frac{Pr^{2}}{2m} + \frac{P\theta}{2m} - V)$$

$$= \partial H(\tau_{1}\theta, \dot{\tau}, \dot{\theta}) = \frac{P_{1}^{2}}{2m} + \frac{P_{0}^{2}}{2m} + V(\tau)$$
Hamilton's equations

For r:

$$\dot{Y} = \frac{\partial H}{\partial P_r} = \frac{P_r}{m}$$
, $\dot{P}_r = -\frac{\partial H}{\partial Y} = \frac{P_0}{mr^3}$, $\frac{\partial V}{\partial r}$
 $\dot{P}_r = \frac{\partial H}{\partial P_0} = \frac{P_0}{mr^2}$, $\dot{P}_{\theta} = \frac{-\partial H}{\partial \theta} = 0 \Rightarrow P_{\theta} = P_{\theta}^{(0)}$
 $\dot{P}_{\theta} = \frac{\partial H}{\partial \theta} = 0 \Rightarrow P_{\theta} = P_{\theta}^{(0)}$
 $\dot{P}_{\theta} = \frac{P_{\theta}}{mr^2}$, $\dot{P}_{\theta} = \frac{-\partial H}{\partial \theta} = 0 \Rightarrow P_{\theta} = P_{\theta}^{(0)}$
 $\dot{P}_{\theta} = \frac{P_{\theta}}{mr^2}$

(3) Simple Pendulum:

The Lagrangian is

$$\mathcal{L} = \frac{1}{2}ml^2\dot{\theta}^2 + mgl\cos\theta$$
, then

$$P_{\theta} = \frac{\partial \mathcal{I}}{\partial \dot{\theta}} = m l^2 \dot{\theta} \implies \dot{\theta} = \frac{P_{\theta}}{m l^2}$$

The Hamiltonian is

$$H = \dot{\Theta} P_{\Theta} - \mathcal{L} = \frac{P_{\Theta}^2}{ml^2} - \left(\frac{1}{2}\frac{P_{\Theta}^2}{ml^2} + \frac{1}{m} \frac{P_{\Theta}^2}{ml^2}\right)$$

$$=) H = \frac{P_0}{2ml^2} - m_g l \cos \theta$$

H's equations

$$\dot{\Theta} = \frac{\partial H}{\partial P_{\Theta}} = \frac{P_{\Theta}}{ml^2}, \qquad \dot{P} = \frac{\partial H}{\partial \Theta} = -mglsin\Theta$$

If
$$\theta$$
 is small $\Rightarrow \cos\theta \approx 1 - \frac{\theta^2}{2}$, then

$$H = \frac{p_{\theta}^2}{2} - mgl\left(1 - \frac{\theta^2}{2}\right) = \frac{p_{\theta}^2}{2ml^2} + \frac{mgl\theta^2}{2} + const$$
(ignore)

0

l

h M

-mgl

This is typical of a Harmonic Oscillator

$$H = \alpha p^2 + \beta \theta^2$$

4.3 Functions on phase space (9, 2)

Calculating

$$\frac{dF}{dt} = \frac{\partial F}{\partial t} + \sum_{i=1}^{N} \left(\dot{q}_i \frac{\partial F}{\partial q_i} + \dot{p}_i \frac{\partial F}{\partial p_i} \right)$$

$$\Rightarrow \frac{dF}{dt} = \frac{\partial F}{\partial t} + \sum_{i=1}^{N} \left(\frac{\partial H \partial F}{\partial p_i \partial q_i} - \frac{\partial H \partial F}{\partial q_i \partial p_i} \right)$$

Suppose
$$\frac{\partial E}{\partial t} = 0$$
 (F does not depend explicitly on t)

$$\frac{dF}{dt} = \sum_{i=1}^{M} (\frac{\partial E}{\partial q_{i}} \frac{\partial H}{\partial p_{i}} - \frac{\partial E}{\partial q_{i}} \frac{\partial H}{\partial q_{i}}) = \frac{4}{4}F_{,H}^{2}$$
Comment: F is conserved when

$$\{F,H\}=0 \qquad (\frac{\partial F}{\partial t}=0)$$

$$\frac{Example:}{Example:} H if \frac{\partial H}{\partial t}=0, \quad \frac{dH}{dt}=\{H,H\}=0 \quad (i.e. F=H)$$
Note: If F and G are 2 functions on phase space then

$$\left\{F,G\right\} = \sum_{i=1}^{N} (\frac{\partial E}{\partial q_{i}} \frac{\partial G}{\partial p_{i}} - \frac{\partial E}{\partial S} \frac{G}{\partial q_{i}})$$
Properties of Poisson brackels
(i) Antisymmetric

$$\left\{F,G\right\} = -\{G,H\}$$
(i) Jacobi identity

$$\left[F,[G,H]\right] = F(GA-HF) = (GA-HF) = 0$$
A bit like matrix commutators

$$\left[F,[G,H]\right] = F(GA-HF) = (GA-GF) + (FA-GF) + (FA-GF$$

 $\Rightarrow \dot{q}_{k} = \sum_{i=1}^{N} \delta_{ki} \frac{\partial H}{\partial P_{i}} = \frac{\partial H}{\partial P_{k}}$

and $\dot{p}_{k} = \{P_{k}, H\}$

Calculate

$$\{q_{k}, q_{k}\} = \sum_{i=1}^{N} \left(\frac{\partial q_{k}}{\partial q_{i}} \frac{\partial q_{l}}{\partial p_{i}} - \frac{\partial q_{k}}{\partial q_{i}} \right) = 0$$

$$i = 1 \qquad = 0 \qquad = 0$$

$$\{P_k, P_k\} = \sum_{i=1}^{N} \left(\frac{\partial P_k}{\partial q_i} \frac{\partial P_k}{\partial p_i} - \frac{\partial P_k}{\partial p_i} \frac{\partial P_k}{\partial q_i}\right) = 0$$

$$\{q_{k}, p_{j}\} = \sum_{i=1}^{N} \frac{\partial q_{k} \partial p_{j}}{\partial q_{i}} - \frac{\partial q_{k} \partial p_{j}}{\partial p_{i}}$$

$$= \sum_{i=1}^{N} \delta_{ki} \delta_{ki} - 0$$